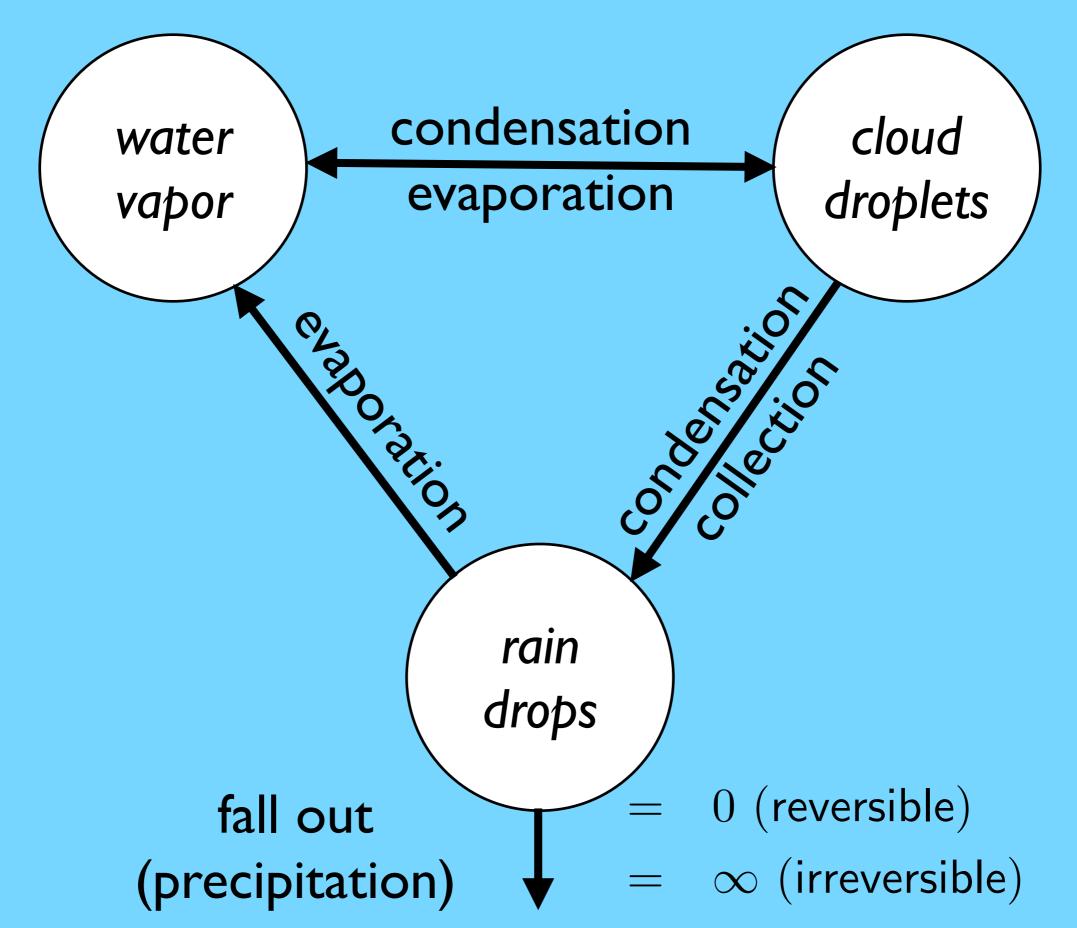
## **Simplified Microphysics**



# **Simplified Microphysics**

$$\frac{d\theta}{dt} = \frac{L}{c_p \bar{\pi}} C$$

$$\frac{dw}{dt} = -C + E_r$$

$$\frac{dl}{dt} = C - A_r$$

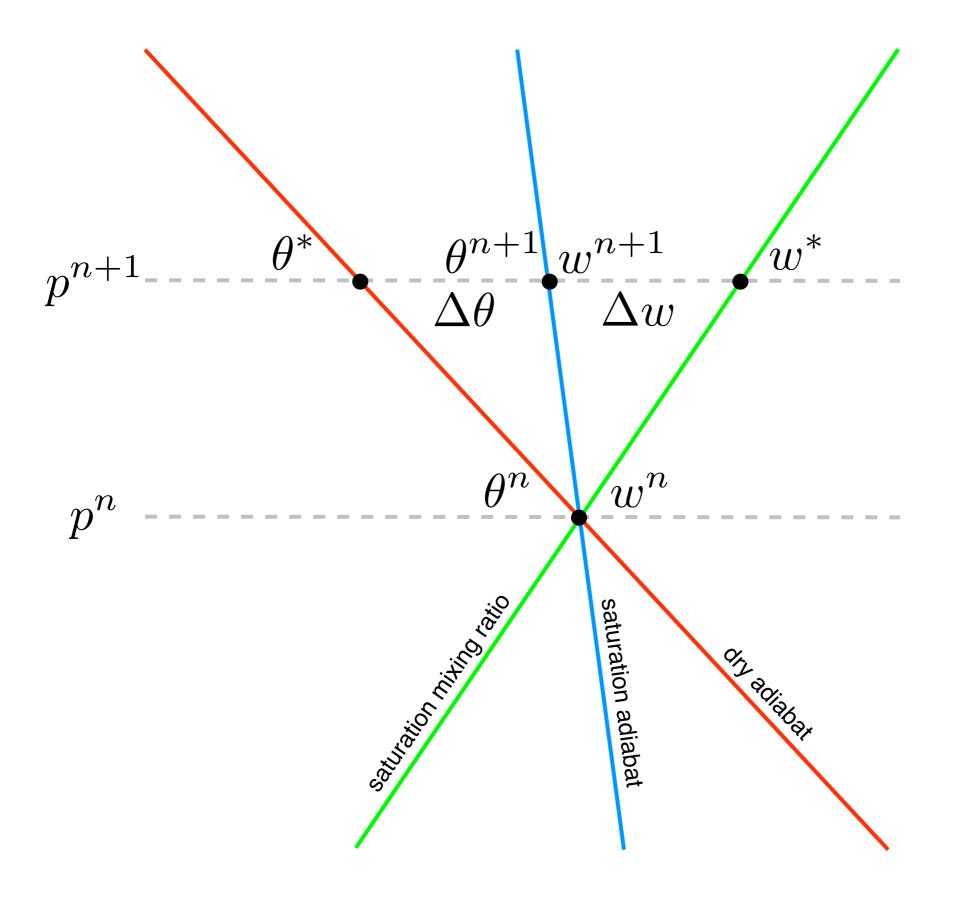
$$\frac{dr}{dt} = P_r + A_r - E_r$$

$$P_r = 0$$
 (reversible)  
 $P_r = \infty$  (irreversible)

# **More Simplified Microphysics**

$$\frac{d\theta}{dt} = \frac{L}{c_p \bar{\pi}} C$$
$$\frac{dw}{dt} = -C$$
$$\frac{dl}{dt} = C - A_r$$

$$A_r = 0$$
 (reversible)  
 $A_r = \infty$  (irreversible)

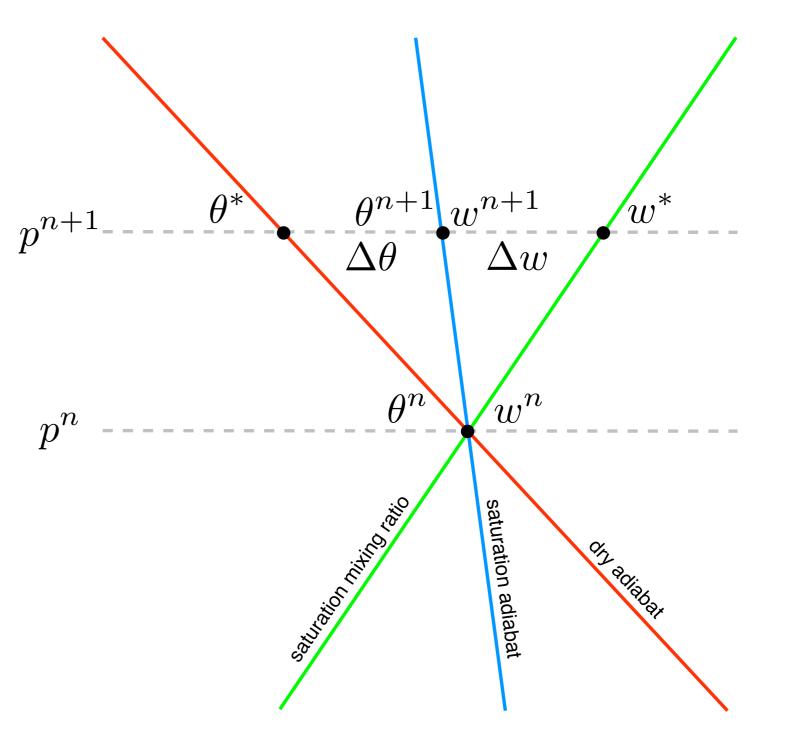


 Adiabatic. No phase changes involving cloud droplets (C=0):

 $\theta^n, w^n \to \theta^*, w^*$ 

2. Isobaric. Only phase changes involving cloud droplets operate (|C| > 0):

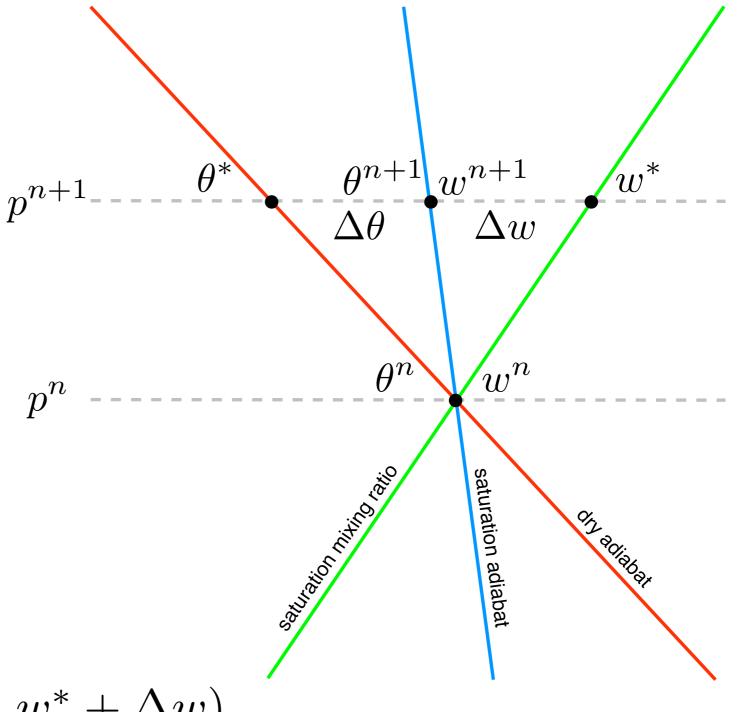
$$\theta^*, w^* \to \theta^{n+1}, w^{n+1}$$

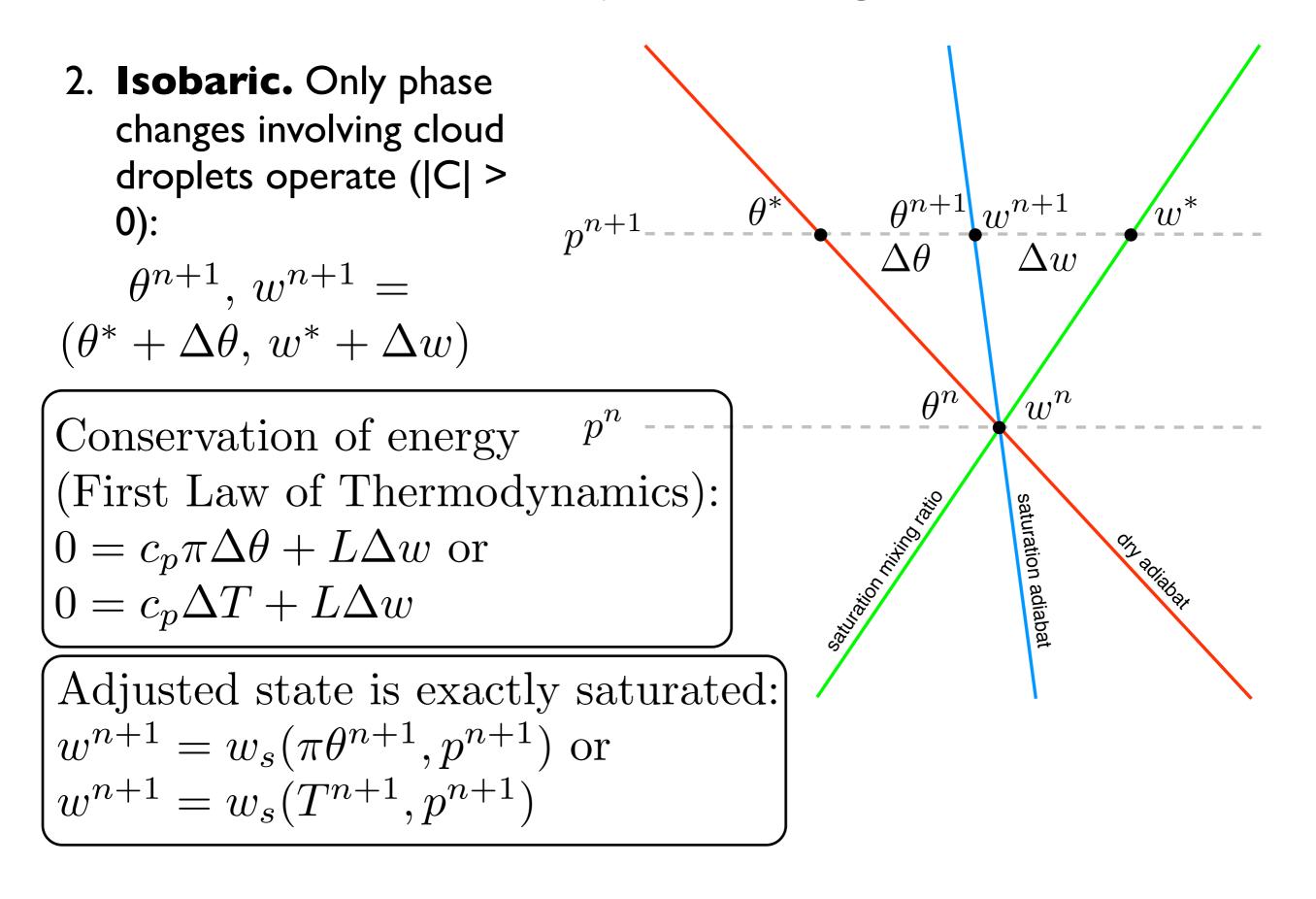


- 1. Adiabatic. No phase changes involving cloud droplets (C=0):  $\theta^n, w^n \rightarrow \theta^*, w^*$  $(\theta^*, w^*) = (\theta^n, w^n)$
- 2. Isobaric. Only phase changes involving cloud droplets operate (|C| > 0):

$$\theta^*, w^* \to \theta^{n+1}, w^{n+1}$$

 $\theta^{n+1}, w^{n+1} = (\theta^* + \Delta\theta, w^* + \Delta w)$ 





$$0 = c_p \pi \Delta \theta + L \Delta w$$
$$w^{n+1} = w_s(\pi \theta^{n+1}, p^{n+1})$$

#### which can be written as

$$0 = c_p \pi (\theta^{n+1} - \theta^*) + L(w^{n+1} - w^*)$$
$$w^{n+1} = w_s(\pi \theta^{n+1}, p^{n+1})$$

$$\theta^{n+1} + \gamma w^{n+1} = \theta^* + \gamma w^* \tag{5}$$

where 
$$\gamma \equiv L/(c_p \bar{\pi})$$
.

conservation of suspended water mixing ratio (vapor and cloud droplets),

$$w^{n+1} + l^{n+1} = w^* + l^*. ag{6}$$

We first assume that the air will be exactly saturated after adjustment, so that

$$w^{n+1} = w_s(T^{n+1}, p^{n+1}), (7)$$

where  $w_s(T, p)$  is the saturation mixing ratio,

$$w_s(T,p) = 0.622 \frac{e_s(T)}{p - e_s(T)},$$
(8)

and  $e_s(T)$  is the saturation vapor pressure. One may use Bolton's (1980) formula for  $e_s(T)$ :

$$e_s(T) = 6.112 \exp\left(\frac{17.67T_c}{T_c + 243.5}\right),\tag{9}$$

where  $e_s$  is in mb,  $T_c = T - T_0$ , and  $T_0 = 273.15$  K.

Equation (7) closes the set (5), (6), and (7). However, this set must be solved iteratively because  $w_s$  is a non-linear function of T. To obtain a direct (non-iterative) solution, expand  $w_s$  in a Taylor series in T about  $w_s(T^*, p^{n+1})$  and neglect all terms of second and higher order:

$$w^{n+1} \approx w_s(T^*, p^{n+1}) + \left(\frac{\partial w_s}{\partial T}\right)_{T=T^*, p=p^{n+1}} (T^{n+1} - T^*).$$
 (10)

The set (5), (6), and (10) can now be solved algebraically for  $\theta^{n+1}$ ,  $w^{n+1}$ , and  $l^{n+1}$ .

To solve the set, we first write (10) in terms of  $\theta$  instead of T:

$$w^{n+1} = w_s^* + \alpha^* (\theta^{n+1} - \theta^*), \tag{11}$$

where  $w_s^* \equiv w_s(T^*, p^{n+1}), \alpha^* \equiv \alpha(T^*, p^{n+1})$ , and

$$\alpha(T,p) \equiv 0.622 \frac{\pi p}{(p-e_s(T))^2} \left(\frac{de_s}{dT}\right)_T.$$
(12)

For  $de_s/dT$ , one may use the Clausius-Clapeyron equation:

$$\frac{de_s}{dT} = \frac{Le_s}{R_v T^2},\tag{13}$$

where  $L = 2.5 \times 10^6$  J/kg and  $R_v = 461.5$  J/(kg K). Now use (11) in (5) to eliminate  $w^{n+1}$ . Then solve for  $\theta^{n+1}$ :

$$\theta^{n+1} = \theta^* + \frac{\gamma}{1 + \gamma \alpha^*} (w^* - w_s^*).$$
(14)

Once  $\theta^{n+1}$  is known from (14), we can immediately obtain  $w^{n+1}$  from (11), and  $l^{n+1}$  from (6).

If  $w^{n+1} \leq w^* + l^*$ , then (6) implies that  $l^{n+1} \geq 0$ . If  $w^{n+1} > w^* + l^*$ , (6) implies that  $l^{n+1} < 0$ , which means that our assumption of saturation is incorrect. Therefore,

$$l^{n+1} = 0$$

replaces (10). Then (5) and (6) become

$$w^{n+1} = w^* + l^*,$$

n + 1 n + (n + 1)