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Quantifying the uncertainty of the annular mode time scale and the role of the stratosphere

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Abstract The proper simulation of the annular mode time scale may be regarded as an important benchmark for climate models. Previous research demonstrated that this time scale is systematically overestimated by climate models. As suggested by the fluctuation-dissipation theorem, this may imply that climate models are overly sensitive to external forcings. Previous research also made it clear that calculating the AM time scale is a slowly converging process, necessitating relatively long time series and casting doubts on the usefulness of the historical reanalysis record to constrain climate models in terms of the annular mode time scale. Here, we use long control simulations with the coupled and uncoupled version of the GFDL climate model, CM2.1 and AM2.1, respectively, to study the effects of internal atmospheric variability and forcing from the lower boundary on the stability of the annular mode time scale. In particular, we ask whether a model's annular mode time scale and dynamical sensitivity can be constrained from the 50-year-long reanalysis record. We find that internal variability attaches large uncertainty to the annular mode time scale when diagnosed from decadal records. Even under the fixed forcing conditions of our long control run at least 100 years of data are required in order to keep the uncertainty in the annular mode time scale of the Northern Hemisphere to 10 %; over the Southern Hemisphere, the required length increases to 200 years. If nature's annular

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² Present Address: School of Earth and Environmental Sciences, Seoul National University, Seoul, South Korea mode time scale over the Northern Hemisphere is similarly variable, there is no guarantee that the historical reanalysis record is a fully representative target for model evaluation. Over the Southern Hemisphere, however, the discrepancies between model and reanalysis are sufficiently large to conclude that the model is unable to reproduce the observed time scale structure correctly. The effects of ocean coupling lead to a considerable increase in time scale and uncertainty in time scale, effects which are noticeable in both troposphere and stratosphere. We further use the model simulation to investigate the dynamical coupling between the stratosphere and the troposphere from the perspective of the annular mode time scale. Over the Northern Hemisphere, we find only weak indication for influences from stratosphere-troposphere coupling on the annular mode time scale. The situation is very different over the Southern Hemisphere, where we find robust connections between the annular mode time scale in the stratosphere and that in the troposphere, confirming and extending earlier results of influences of stratospheric variability on the troposphere.

Keywords Annular mode time scale · Internal variability · Stratosphere–troposphere coupling · Dynamical sensitivity

1 Introduction

Climate sensitivity is a remarkably important characteristic of climate because many aspects of climate change scale almost linearly with it (Meehl et al. 2007). Research on climate sensitivity has now been ongoing for decades, but despite this the estimated range of climate sensitivity and the associated uncertainty essentially remained unchanged over this period (Knutti et al. 2006; Houghton et al. 2001;

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Randall et al. 2007). The main reasons for this lack of progress are the shortage of reliable observations and uncertainties in the formulation of climate models.

Recently, it has been suggested that the persistence time scale of major modes of extratropical variability in the atmosphere, also known as the annular modes (Thompson and Wallace 2001), could provide another important measure of sensitivity (Gerber et al. 2008a; Ring and Plumb 2008; Chen and Plumb 2009). This measure has been termed "dynamical sensitivity" to emphasize that it expresses the sensitivity of the circulation to climate change (Grise and Polvani 2014). As predicted by the fluctuation-dissipation theorem (Leith 1975), the equilibrium response to external forcings and thus dynamical sensitivity should be proportional to the persistence time scale of the annular modes (hereafter simply AM time scale or τ). Thus, comparing the time scale between models and observations could provide a useful alternative for understanding how realistic a model's dynamical sensitivity is.

Previous studies already investigated the AM time scale from observational (Baldwin et al. 2003) and modeling (Gerber et al. 2008a, b, 2010; Son et al. 2008) data. Gerber et al. (2008b, 2010) found that the AM time scale is systematically overestimated in climate models, particularly during summer over the Southern Hemisphere (SH). Likewise, the AM time scale seems to be unrealistically long in more simple models (Gerber et al. 2008a).

In order to test dynamical sensitivity using the AM time scale one needs to determine the time scale robustly. However, the AM time scale is highly sensitive to internal atmospheric variability, complicating reliable estimates. For example, Chan and Plumb (2009) find that in idealized models irregular and unpredictable regime shifts of the jet stream interfere with the determination of the AM time scale. Further, Gerber et al. (2008a) suggest that about 30 years of perpetual January model simulations are required to determine a time scale of 25 days at a 10 % accuracy. However, their measure is only an approximation and does not take into account complicating effects from the annual cycle and from variations in external forcings in more realistic data.

The difficulty in determining τ becomes evident from Fig. 1, showing the reanalysis derived time scale as a function of season and height for two 25-year-long non-overlapping periods. Although the two resulting τ structures are similar, there exist also important differences that cast doubts on the interpretation of the outcomes. For example, the result from the first period suggests that the wintertime peak in τ occurs first in the stratosphere and then in the troposphere, but the second period shows the opposite behavior. A similar analysis conducted by Baldwin et al. (2003) using the same data but for one longer period (1958–2002) suggests that the stratospheric peak in τ



Fig. 1 NAM time scale (in days) as a function of month of the year and pressure, derived from the first (1959–1983) and last 25 years (1984–2008) of the NCEP/NCAR reanalysis

precedes that in the troposphere, similar to what is shown in Fig. 1a for the first period. Some of these discrepancies might be related to inconsistencies in the reanalysis or trends associated with climate change. An equally valid explanation, however, is that the time scale is highly sensitive to natural variability in the underlying data and that the process of calculating the time scale converges only slowly to its actual value. In other words, 25 or 50 years of observations may not be sufficient to derive a reliable estimate of τ .

The goal of the present study is to investigate the AM time scale in a model and to understand the role of internal atmospheric variability for determining the time scale. We try to answer the practical question of how many years of data are actually required to reduce the uncertainty in τ below a certain level. A related question is to what extent differences between models and reanalysis are indicative for real model biases and to what extent they are simply due to internal atmospheric variability, also known as noise. We address our questions from long control simulations with coupled and uncoupled climate models, which because of their long length allows us to answer these questions with a high degree of certainty.

Examining the differences between the coupled and uncoupled model enables us to distinguish between the influences from internal atmospheric dynamics and that from persistent lower boundary condition forcing on the annular mode time scale and its uncertainty. To the extent that our models are representative for other models, our finding then provides new insight into the shortcoming of models in reproducing the observed time scale. As we will show, the limited length of the historical observations combined with internal atmospheric variability leads to large uncertainty in determining the AM time scale. Over the Northern Hemisphere (NH), the reanalysis derived time scale therefore does not provide a useful constraint for dynamical sensitivity. The situation over the SH is somewhat different in that we find that the model has a significant bias in reproducing the observed time scale structure.

Another question to be addressed in our study is whether the stratosphere has as a detectable influence on the troposphere in terms of τ . Such an influence might be related to a dynamical downward influence from the stratosphere into the troposphere. Based on the observation that the winter time peak in the tropospheric AM time scale over the Northern Hemisphere is preceded by a similar peak in the stratosphere, Baldwin et al. (2003) were the first to suggest that such an influence exists. More recently, Simpson et al. (2011) found that indeed a dynamically active stratosphere serves to lengthen the time scale of the tropospheric annular mode during winter. In the present study we come to a similar conclusion, but we find that the strength of the coupling in terms of the time scale is weaker over the NH than over the SH.

The outline of this paper is as follows. In Sect. 2, we describe our data and methodology. Section 3 compares the observed AM time scale with that derived from the model simulation. In Sect. 4, we investigate the uncertainty of the AM time scale when the calculations are based on the 50-year-long historical record. Section 5 explores the extent to which agreement between model and observations based on 50 years of data can be used to evaluate a model. In Sect. 6 we investigate the connection between the stratosphere and troposphere in terms of the time scale and find that there are important differences between the NH and the SH. In Sect. 7, we investigate the uncertainty of the AM time scale structure as a function of length of the underlying data, which we hope provides useful information for other similar studies. Our final section offers a summary and interpretation of our results.

2 Data and methodology

2.1 Data

We use daily zonal mean geopotential height fields poleward of 20° from the National Centers for Environmental Prediction/National Center for Atmospheric Research (NCEP/NCAR) reanalyses (Kalnay et al. 1996), which are widely employed in annular mode studies (Baldwin and Thompson 2009; Thompson and Wallace 2000). The reanalysis dataset is available at 17 vertical levels from 1000 to 10 hPa, and our calculations are based on the 50-yearlong period from 1959 to 2008.

The model simulated annular mode data are derived from long control simulations performed with two models. First, we utilize a 4000-year-long preindustrial control run with the coupled climate model CM2.1, developed at the Geophysical Fluid Dynamics Laboratory (GFDL Global Atmospheric Model Development Team 2004; Delworth et al. 2006). Second, we utilize a 2000-year long control integration with the uncoupled atmosphere-only companion model of CM2.1, also known as AM2.1. The prescribed SST and sea ice distribution (hereafter simply SST) for the uncoupled model were derived from averaging the coupled model's SST over many years (Staten et al. 2014). All external forcings (greenhouse gases, ozone, solar influences, and aerosol) in both simulations were held fixed at preindustrial levels. The spatial resolution of both models in the atmosphere is 2 degree latitude by 2.5 degree lon-

gitude with 24 vertical levels up to 10 hPa. The models produce realistic climates (Reichler and Kim 2008), and despite their relatively low vertical resolution, they simulate the annular mode variability and its coupling between the stratosphere and troposphere quite well (Reichler et al. 2012).

2.2 Methodology

Our procedure to compute the annular mode indices follows the method proposed by Baldwin and Thompson (2009). Specifically, it is based on the leading empirical orthogonal function (EOF) of daily zonal mean geopotential height anomaly fields at each pressure level, with anomalies being defined as deviations from the daily varying climatology with the global mean removed. The annular mode index is the principal component time series of the leading EOF. Similar to previous studies (Gerber et al. 2008b; Baldwin et al. 2003), the AM time scale is defined as the e-folding time of the autocorrelation function of the annular mode index. The calculation of the autocorrelation function is based on the daily annular mode index time series for all years. To calculate the autocorrelation for a given day of the year, we multiply each year of the original annular mode time series with a Gaussian kernel with a full width at half maximum of 60 days centered on the day under consideration. The result is a time series with yearly oscillating amplitudes, maximum amplitudes for dates close to the central date, and decreasing amplitudes for dates distant to it. We next create 180 similar time series, but at lags for ± 1 to ± 90 days with respect to the central date, and then correlate each with the time series for the central date. Since the resulting autocorrelation function is far from exponential (Ambaum and Hoskins 2002) and since the e-folding time scale is only meaningful if the autocorrelation function drops off exponentially, we base the calculation of the time scale off an exponential least-square fit to the empirically

determined autocorrelation function. The AM time scale is then given by the day the fitted autocorrelation function drops to a value of 1/e. We repeat these calculations for each day of the year.

The time scale from the models is calculated by splitting the 4000 and 2000-year-long simulations into several N-year-long segments (denoted as model segments). In most cases, we use N = 50, which allows a nice comparison with the reanalysis derived time scale and which results in 80 and 40, respectively, different segments. We calculate the AM time scale individually for each segment and then use the mean time scale from all segments as our best estimate. The variability of the AM time scale is derived from the standard deviation across all segments, and our measure of uncertainty is given by the ratio between the standard deviation and the mean.

3 Comparison between model and reanalysis

We begin with a comparison of reanalysis and model derived AM time scale. This discussion is based on Fig. 2, which shows the time scale structures of the Northern and Southern Annular Modes (NAM and SAM) derived from 50 years of reanalysis data and 4000 and 2000, respectively, years of model data. We first note that the AM time

scale calculated from zonal mean fields is very similar to that from using two-dimensional data (e.g., Figure 1 of Baldwin et al. 2003). Key features of the reanalysis-derived time scale structure are well captured by the models. This includes tropospheric peaks in the NAM during boreal winter and in the SAM during austral spring. Further, in the stratosphere both reanalysis and models exhibit much longer time scales than in the troposphere, with two distinct maxima at 10 hPa in summer and at 100 hPa in winter for the NAM, and a time scale that is long throughout the year for the SAM.

The model derived time scales also exhibit important differences with respect to the reanalyses. First, the tropospheric peaks both for NAM and SAM are too broad in the models, a deficiency common to many models (Gerber et al. 2008b, 2010). Second, the tropospheric peak of the NAM in the models occurs in February, which is about 1–2 months delayed with respect to the reanalyses. Third, the models generally underestimate the stratospheric AM time scales over both hemispheres, whereas they overestimate the tropospheric time scale, particularly for the coupled model and in the SH. As predicted by the fluctuation–dissipation theorem, the linear relationship between AM time scale and dynamical sensitivity suggests that the model's sensitivity is too large in the troposphere and probably too small in the stratosphere to the extent that these

Fig. 2 NAM and SAM time scale (in days) for (top) NCEP/ NCAR reanalysis (1959–2008), (middle) the coupled (CM, mean over eighty 50-year-long segments) and uncoupled (AM, mean over forty 50-year-long segments) models, and (bottom) the percentage difference between the two models (CM-AM). The time scale at 1000 hPa is derived from zonal mean sea level pressure; for all other levels, zonal mean geopotential heights are used



biases are related to internal atmospheric dynamics and not to external persistent forcings.

We now discuss in more detail the differences in time scale between the coupled and uncoupled simulations (Fig. 2g, h). The prevalence of reddish colors in the difference plots indicates that, on average, the coupled model has a longer time scale than the uncoupled model. This effect is unsurprising as low-frequency variations of the ocean increase the atmosphere's interannual variability and hence its time scale (Keeley et al. 2009). In addition, it is well known that prescribed SSTs induce artificially large turbulent heat fluxes at the atmosphere's lower boundary, reduce atmospheric variability (Barsugli and Battisti 1998), and hence decrease atmospheric persistence. A third factor which may play some role is that the climatological base state between the two models is not exactly the same; this may cause subtle changes in the internal dynamics and processes that determine the time scale (e.g., positive dynamical feedbacks).

For the NAM (Fig. 2g), ocean coupling leads to a modest ~20 % increases in tropospheric τ (from ~9 to ~12 days) during winter and spring. Generally, this is the time when atmosphere–ocean interaction is strongest and the El Nino Southern Oscillation is most active. However, the exact reason why the coupling effect on τ is most pronounced during April is not clear to us. It may be related to the winterto-summer transition of the atmospheric circulation, the timing of which is likely to be more variable in the coupled model. For the SAM (Fig. 2h), the increase in τ from ocean coupling is more robust (ca. +50 %) than for the NAM, it affects the entire atmospheric column, and is most pronounced during SH summer (DJF). The peak influence during summer may be related to enhanced variability in the timing of the breakdown of the SH polar vortex, which is believed to be responsible for the overall increase in time scale during this time of the year.

4 Time scale uncertainty in the historical record

We next examine the uncertainties associated with the AM time scale when 50-year-long data such as the NCEP/ NCAR reanalysis are available. To this end we split the 4000-year-long (2000-year-long) coupled (uncoupled) model simulation into eighty (forty) different 50-year-long segments and calculate the time scale individually for each segment. The standard deviation (or variability) amongst the eighty coupled model segments exhibits a similar structure as the time scale itself. The variability of the NAM time scale of the coupled model (Fig. 3a) displays distinct maxima of 4–5 days at 100 hPa in early winter and at 10 hPa in summer. In the lower stratosphere and in the troposphere, this large variability persists out to spring. A similar variability in the seasonal timing of the peak time scale of the NAM amongst individual 50 year selections of the same simulation was found by Simpson et al. (2011). The variability of the SAM time scale (Fig. 3b) is generally much larger than that of the NAM, somewhat consistent with the larger magnitude of the SAM time scale itself (Fig. 2d). In the lower stratosphere and also in the troposphere, the SAM time scale is most variable during late spring and early summer, exceeding values of 15 days.

Fig. 3 Variability and uncertainty of model derived time scale. Shown are (top) standard deviation of the time scale (in days), derived from 80 samples of 50-year-long segments from the coupled model (CM), (middle) uncertainty of the time scale of the coupled model (in %), given by the standard deviation divided by the mean uncertainty, and (bottom) the difference in uncertainty (in %) between the coupled and uncoupled model. Using F-statistics and neglecting differences in mean timescale between the two models, differences in uncertainty of about 20 % are statistically significant



Fig. 4 Coupled model derived time scale for selected 50-yearlong segments. The *top* (*bottom*) *panels* are examples that are in good (*poor*) agreement with the reanalysis. *Numbers at the top* right of each *panel* indicate the pattern correlation between reanalysis (1959–2008) and model derived time scale structure



The large variability of the NAM time scale in the lower stratosphere during winter is likely related to different forms of dynamical variability in the stratosphere, which includes strong vortex events as well as stratospheric sudden warming (SSW)-like events. SSWs occur in winter and represent abrupt weakenings or even complete break-downs of the polar vortex, creating pronounced NAM variability in the lower stratosphere and also in the troposphere (Simpson et al. 2011; Baldwin et al. 2003). The variable timing and long persistence of weak and strong vortex events and also their infrequent occurrence likely contribute to the wintertime maximum of the NAM time scale in the stratosphere and troposphere (Fig. 2c). This and also interdecadal intermittency in the occurrence of SSWs (Schimanke et al. 2011; Reichler et al. 2012) help to explain the large wintertime variability of the NAM time scale (Fig. 3a). Over the SH there are no SSWs. Instead, during spring there is large variability in the exact timing of the annual breakdown of the polar vortex, which is an important factor for interannual variability in the SAM during this time of the year. This explains the increase in the SAM time scale (Fig. 2d) and its variability (Fig. 3b) during spring.

A more objective measure for the uncertainty of the AM time scale is derived by taking the ratio between the standard deviation and the mean of the time scale (Fig. 3c, d). This eliminates the effect of the mean on the variability. In most cases the annular mode uncertainty derived from 50-year-long segments amounts to 10-20 % of the mean, but it occasionally exceeds 25 %. We note that the uncertainty structure of the NAM (Fig. 3c) closely resembles that of the standard deviation (Fig. 3a), whereas the SAM (Fig. 3b, d) does not show this effect.

The lower panels of Fig. 3 show the percentage difference in uncertainty between the coupled and uncoupled model. The reddish colors indicate that ocean coupling leads to a general increase in uncertainty, as one might expect from the added atmospheric variance from the lower boundary condition forcings. However, the structure of uncertainty increase does not necessarily resemble the structure in time scale increase (Fig. 2g, h). In particular for the NAM there are increases in uncertainty (up to 60 %) during most times of the year and throughout the entire atmosphere, which reflects intense vertical coupling between the surface, troposphere, and stratosphere. The situation for the SAM is somewhat different, with relatively modest increases in uncertainty concentrated in the stratosphere and during SH spring. The weakness of the ocean coupling effect over the SH is perhaps related to the fact that the southern polar cap is entirely covered by land.

5 Comparison of model and reanalysis derived AM time scales

We now explore to what extent the differences between reanalysis and model derived time scale (Fig. 2) are systematic and indicative for actual model deficiencies and to what extent they are related to internal atmospheric variability and the difficulty to calculate a stable τ estimate from a record that is only 50 years-long. We begin by selecting, from the 80 segments, examples that show a particularly good or bad agreement with the reanalysis derived time scale structure (Fig. 4). Our measure of goodness is based on the pattern correlation between the observed and model derived time scale structure for all levels and months of the year. The two good examples (Fig. 4, top panels) are in reasonable agreement with the observations (Fig. 2a, b) and are certainly more realistic than the mean time scale derived from all segments (Fig. 2c, d). This is true with respect to the magnitude, width, and timing of the stratospheric and tropospheric peaks. The two examples of particularly poor agreement (bottom panels Fig. 4) hardly capture the observed seasonal cycle of the time scale, with peaks that have different timing, multiple maxima, no maxima, or are much too broad in structure. These examples demonstrate that internal variability has a huge effect on the outcomes of the time scale structure and that one must be cautious when judging models based on their performance in reproducing the observed time scale of the 50-year-long reanalysis record.

Kidston and Gerber (2010) showed that the time scale in CMIP3 models is correlated with biases in the latitude of the jet stream. This motivates an investigation of whether a similar connection between climatological biases and annular mode time scale exists amongst the segments of our single coupled model (Fig. 5). As potential predictors for variations in time scale we utilize, individually for each 50-year-long segment, the climatological zonal means of (Fig. 5a, b) the maximum value of the surface wind, corresponding to the strength of the eddy driven jet (Fig. 5c, d) the latitude location of the surface wind maximum, corresponding to the position of the eddy driven jet (Fig. 5e, f) the maximum value of the winds at 10 hPa, corresponding to the strength of the polar vortex, and (Fig. 5g, h) the latitude of the wind maximum at 10 hPa, corresponding to the position of the polar vortex. We present the outcomes for the (left panels) NAM during DJF and (right panels) SAM during SON, noting that the results for other seasons are similar. Figure 5 exhibits a large scatter in time scale amongst the individual segments, indicative for the time scale uncertainty, but none of the four selected predictors forms robust correlations with that scatter. The only exemption is perhaps the magnitude of the maximum surface wind (Fig. 5a, b), which for both the NAM and SAM exhibit a weakly negative linear relationship with the tropospheric time scale. Likewise, there is some modest positive correlation for the latitude of the surface wind maximum and the time scale of the tropospheric NAM (Fig. 5c). However, this positive correlation contradicts Kidston and Gerber (2010), who find that an equatorward shifted SH jet leads to enhanced SAM persistence. We conclude that climatological wind shifts of shorter model segments are mostly unrelated to variations in time scale, that such shifts are mostly due to internal variability, and that therefore longer than 50-year-long integrations are needed to use the annular mode time scale for model assessment.

We next investigate the similarity of the τ structure between model and reanalysis more systematically. The top panels in Fig. 6 show scatter plots of root-mean-square errors (RMSEs) in τ structure between individual segments and reanalysis, separately for the stratosphere and the troposphere. Pattern correlations for the entire domain are also shown in color. In calculating the RMSEs we consider



Fig. 5 Time scale as a function of wind biases. *Red dots* show (*red*) time scale of eighty coupled model segments as a function of their corresponding climatological zonal mean of (a, b) surface wind maximum, (c, d) latitude of the wind maximum, (e, f) wind maximum at 10 hPa, and (g, h) latitude of the wind maximum at 10 hPa. *Black dots* in each panel correspond to the NNR reanalysis

levels from 200 to 30 hPa for the stratosphere and from 1000 to 500 hPa for the troposphere, giving equal weights to each level. Also, we only include the 8 months from September to April because outside this period the time scale exhibits artificially large values in the stratosphere when the annular modes are inactive (Gerber et al. 2010).

For the NAM (Fig. 6a), the RMSE values scatter around 4 days and are about the same in the stratosphere and troposphere. The pattern correlations vary between 0.9 and 0.5.



Fig. 6 Relationship between tropospheric (1000–500 hPa, September–April) and stratospheric (200–30 hPa, September–April) time scale. (*Top*) root-mean-square errors (RMSE) in time scale structure between 80 coupled model derived segments and the reanalysis. (*Bottom*) RMSE between all paired combinations of 80 model derived segments; shown are the outcomes of all 3160 (= $80 \times 79 \div$ 2) unique pairs. In all panels, *color* denotes pattern correlations calculated over all levels (1000–10 hPa) and months (January–December). *Ellipses* are centered on the mean, oriented along the direction of maximum scatter; the half axes of the ellipses measure two standard deviations along the major and minor directions. *Lines* represent the diagonal where tropospheric and stratospheric RMSE match. *Numbers at the bottom* of *each panel* denote correlations between stratospheric and tropospheric scatters, with *numbers in parenthesis* derived from the uncoupled model

Several model segments have high correlations and relatively low RMSEs, suggesting that the well performing example from Fig. 4a is not just due to exceptionally good luck but rather that the model is actually capable of simulating the observed time scale if the conditions related to internal variability are right. Closer inspection reveals that, to first order, a good skill in simulating the observed time scale structure results when the seasonal timing of the stratospheric and tropospheric peaks in AM time scale is well reproduced. Since for the NAM these peaks owe their existence to variability of polar vortex strength (Simpson et al. 2011), which essentially is random in both frequency and seasonal timing (Schimanke et al. 2011; Reichler et al. 2012), we conclude that the scatter of outcomes displayed in Fig. 6a mostly reflects internal variability. In other words, the differences between simulated and observed time scale structure in the NAM are not indicative of systematic model errors-at least not in the model of this study.

The skill of the model in reproducing the observed time scale structure of the SAM is considerably smaller than that for the NAM. The RMSE values for the SAM (Fig. 6b) scatter around 10-20 days with pattern correlations of about 0.6. One important reason for the smaller skill is the inability of the model to reproduce the observed structure of the tropospheric peak at the beginning of December. For example, comparing Fig. 4b with Fig. 2b shows that even the best-selected model segment simulates the observed tropospheric peak in SAM time scale quite poorly. All 80 segments exhibit a similar behavior, with a troposphere peak that is much too broad, too late, and too intense. We therefore conclude that this model exhibits a systematic bias in terms of the SAM time scale structure and that this bias cannot be explained from sampling uncertainty and internal variability. Previous studies already identified this as a typical GCM problem (Gerber et al. 2008b, 2010; Simpson et al. 2011). This is related to the too-late breakdown of the polar vortex (Simpson et al. 2013a) and to a bias in the feedback by planetary scale waves on the SAM (Simpson et al. 2013b).

We next repeat our analysis but using a perfect model approach under which each segment is validated against all other segments of this model. The bottom panels of Fig. 6 show the RMSE values and correlations of the AM time scales among all unique pairs of model segments. For the NAM, the shape of the scatter and also the associated spread of the correlations are close to that of the model-reanalysis validation (Fig. 6a), confirming our interpretation that apparent model errors for the NAM are mostly due to internal variability. For the SAM, however, the skill under the perfect model approach is much higher compared to the model-reanalysis validation. For example, the dominance of reddish colors indicates that the correlations are uniformly high across almost all model pairs, and the RMSE values range only between 5 and 10 days. This is in stark contrast to Fig. 6b for the model-reanalysis comparison, and substantiates our assumption that the model simulated time scale of the SAM is systematically biased with respect to the real atmosphere.

6 Linkages between stratosphere and troposphere

As seen before (Fig. 2), both NAM and SAM time scales exhibit a pronounced seasonal structure, with the lower stratospheric peak in τ being followed by a similar peak in tropospheric τ . This observation led Baldwin et al. (2003) to argue that the tropospheric peak is related to a dynamical influence from the stratosphere on the troposphere. Gerber and Polvani (2009) used an idealized general circulation model to show that realistic stratospheric variability increases the time scale of the tropospheric annular mode, confirming the earlier hypothesis from Baldwin et al. (2003). Similar conclusions were reached by Gerber et al. (2010), using a suite of stratosphere-resolving chemistry climate models, and by Simpson et al. (2011), using the Canadian Middle Atmosphere Model. More precisely, Simpson et al. (2011) found that stratospheric variability approximately doubles the peak of the tropospheric SAM time scale, but that influences other than that from the stratosphere also contribute to the long tropospheric SAM time scale. Here, we further address the issue of stratospheric influences on tropospheric time scale using our large ensembles with the coupled and uncoupled GFDL models, applying a different technique, and extending the analysis also to the NH.

We begin by investigating the relationship between stratospheric and tropospheric RMSE values from individual model segments of the coupled model when one validates against the reanalysis (top panels of Fig. 6). For the NAM (Fig. 6a), there is no obvious relationship between stratospheric and tropospheric RMSE values as seen from the zero correlation of the scatter. For the SAM (Fig. 6b), however, the two seem to be somewhat related, as shown by the modestly positive correlation (r = 0.38) of the scatter and the elliptical shape of its envelope. In other words, for an individual 50-year-long segment the time scale in the troposphere is generally better simulated when the time scale in the stratosphere is also well simulated, and vice versa. The connection between tropospheric and stratospheric time scale becomes even clearer when a perfect model approach is used to validate the model against itself (Fig. 6 bottom). As before, the correlations of the scatter are clearly positive for the SAM (r = 0.59), and even for the NAM they are now somewhat positive (r = 0.14). This indicates that indeed stratospheric and tropospheric time scales are connected, and that in this model the connection is stronger for the SAM than for the NAM. The stronger SAM connection is consistent with the longer time scale of the tropospheric SAM as compared to that of the NAM, which according to the fluctuation-dissipation theorem should be associated with a stronger response to a given stratospheric forcing.

We further investigate the possible influence of the stratospheric time scale on the tropospheric time scale by focusing separately on the timing and on the strength of the peaks in τ (Fig. 7). To this end, we determine for all 50-year-long model segments the date and the strength of maximum τ in the lower troposphere (1000–500 hPa) and in the lower stratosphere (200–30 hPa). For the NAM, the stratospheric and tropospheric peaks occur between November and March (Fig. 7a). The shape of the ellipses of maximum scatter and also the correlation of the scatter itself show that the dates of the tropospheric peaks are essentially unrelated to the dates of the stratospheric peaks (r = 0.08). Also, most scatter symbols are located above the diagonal



Fig. 7 Relationship between stratospheric and tropospheric time scale maxima derived from 80 coupled model segments. a, b Date of lower stratospheric (200-30 hPa) and lower tropospheric (1000-500 hPa) time scale maxima (results are not sensitive to the exact choice of the vertical levels). c, d Relative strength of time scale maxima, given by the ratio of time scale maxima and the mean of all maxima (23 days for stratospheric NAM, 15 for tropospheric NAM, 71 for stratospheric SAM, and 47 for tropospheric SAM). Colored circles show outcomes from individual model segments and black circles indicate the mean. Ellipses are centered on the mean, oriented along the direction of maximum scatter, with the two axes showing four standard deviations along the major and minor direction. Numbers at the bottom are correlations between stratospheric and tropospheric scatters. Numbers in parenthesis indicate correlations derived from the uncoupled model. Black diamonds show outcomes from NCEP/ NCAR reanalysis (1959-2008). Lines in top panels represent matching stratospheric and tropospheric date

line, indicating that the tropospheric peak dates usually lag the stratospheric peaks. Taking the mean over all segments (black circle in the center of the ellipses) we find that in the stratosphere the peak usually occurs during mid-January and in the troposphere at the beginning of March, which is a lag of several weeks. In contrast to the simulations, the reanalyses (black diamond) exhibit about a 1-week delay from the stratosphere into the troposphere and a peak date that occurs about 1 month earlier. Examining the magnitude of the peaks in τ for the NAM (Fig. 7c), we again find only a very weak connection between stratosphere and troposphere (r = 0.09). We note that, despite the differences in the mean timing, the outcomes for the segments as a whole are consistent with the reanalysis since the black diamond symbols for the reanalyses are located well within the uncertainty ellipses of the model. In other words, various segments have

a NAM time scale structure that is similar to the reanalysis, both in terms of peak date and peak magnitude.

For the SAM there is somewhat less scatter in peak dates (Fig. 7b), but as for the NAM the peaks in the stratosphere mostly lead that in the troposphere, which may be somewhat indicative for a stratospheric influence on the troposphere. The SAM peak dates are centered on late November in the stratosphere and on December in the troposphere. Comparing to the NAM and taking into account the seasonal shift of 6 months between the two hemispheres, the SAM dates are shifted late by about 4 months. This shift, which was also noted by Gerber et al. (2008b), is consistent with the fact that variability of the NAM is dominated by vacillations of the polar vortex during mid-winter whereas that of the SAM is related to the changing breakdown date of the polar vortex during spring. For the SAM the stratospheric peak timing is somewhat correlated with the tropospheric peak (r = 0.20). The narrow range in timing and the moderately positive correlation between stratospheric and tropospheric peak timing is likely related to the smaller internal variability and non-existence of SSWs over the SH. As shown in Fig. 7d, there is much more scatter in the magnitude of the peak time scale for the SAM than for the NAM and a quite significant correlation between the stratosphere and the troposphere (r = 0.72). The model outcomes are consistent with the reanalyses (black triangles) for the seasonal timing of the peak. This, however, is not true for the intensity of the peak time scale, which is significantly stronger in the model than for the reanalysis.

The above analysis (Figs. 6, 7) suggests that stratospheric and tropospheric time scale are somewhat connected, in particular over the SH. However, this analysis is based on coupled model data, and it is still unclear to what extent the connection is due to oceanic variability influencing stratosphere and troposphere in similar ways and to what extent it is due to internal atmospheric processes. In order to shed somewhat more light on this question, we repeat our analysis using data from the uncoupled model, where oceanic variability cannot be the source or recipient of influence. The outcomes for the uncoupled model data (see numbers in parenthesis) are similar to the coupled model, except that the stratospheretroposphere relationships become weaker for the SAM and stronger for the NAM. For example, the relatively tight correlation of r = 0.72 in SAM peak time scale reduces to r = 0.32in the uncoupled model. Such a reduction is expected as ocean signals forcing the troposphere, stratosphere, or both are eliminated, and tropospheric variability is damped. In terms of peak time scale (Fig. 7c, d), the remaining relationship between stratosphere and troposphere is weakly positive $(r \sim 0.3)$ and about the same for NAM and SAM. The small correlation is somewhat indicative for a stratospheric influence in lengthening the tropospheric time scale.

7 Uncertainty in time scale

Calculating the time scale of the annular mode is a slowly converging procedure, and multi-decade-long data are needed to arrive at reasonable estimates. From Fig. 3 one can see that the uncertainty is ~15–25 % if the calculations are based on a 50-year-long record. This raises the more general question of how long of a record is required in order to limit the uncertainty of the time scale calculations to a certain value. Our long control simulations provide an excellent opportunity to answer this question by deriving empirical relationships between time scale uncertainty and length of the simulation period.

We employ a procedure similar to that of Fig. 3 and calculate a measure of uncertainty of the AM time scale by determining the standard deviation of the τ structure amongst multiple N-year-long model segments, averaged over all levels and all days of the year. We repeat the calculation for increasing N = (10, 20, 50, 100, 200, 500 years), noting that the total number of segments decreases as N increases. In Fig. 8, the outcomes for the coupled and uncoupled model are shown by the red and blue symbols, respectively.

For a segment length of N = 10 years, the uncertainty is very large and amounts to 40-50 % of the mean time scale. As expected, the uncertainty and its range become smaller as N increases. In order to limit the uncertainty to 10 % at least N = 100 years of data are needed. The lines in Fig. 8 are extrapolations from the empirically determined values for N = 10 years, assuming that uncertainty for longer N scales inversely proportional to the square root of N. Such a scaling is intuitive if one assumes that calculating τ over increasing N is similar to taking the mean τ from multiple (=N/10) 10-year-long segments and knowing that the standard error of the mean equals the sample standard deviation divided by the square root of the sample size. One can see that the extrapolated values are very similar to the empirically determined ones. Therefore, it is reasonable to assume that the extrapolated uncertainties for the reanalysis (black curve) are good approximations for the actual (unknown) uncertainties.

Figure 8 also indicates that the uncertainty of the SAM is larger than that of the NAM (50 vs. 40 % for N = 10). Since our measure of uncertainty is relative to the mean, the increase is indicative for differences between NAM and SAM time scale that go beyond the simple effects of the mean.

Comparing the differences in NAM uncertainty from the coupled (red) and uncoupled (blue) model one finds that, as expected from Fig. 3e, the uncertainty of the uncoupled model is consistently lower, but the differences with respect to the coupled model are not very large. For the SAM, the



Fig. 8 Uncertainty of NAM and SAM time scale structure (averages over all levels and days of the year) as a function of length of the underlying time series for (*black*) the NCEP/NCAR reanalysis, (*red*) the coupled model, and (*blue*) the uncoupled model, all calculated from multiple segments of given length N. *Grey symbols* indicate outcome for NCEP/NCAR reanalysis from using bootstrapping (100 times, with replacement) data within the N-year-long segments.

Uncertainty is defined as in Fig. 3. *Circles* are actual calculations (slightly shifted along the *x*-axis for clarity), and curves represent extrapolations from N = 10 years using the analytical expression (inversely proportional to the square root of the length of the segment, see text). *Error bars* denote ± 2 standard deviations from the mean, calculated from bootstrapping by randomly selecting a subset of five segments with replacement and repeating this 100 times





Fig. 9 Multiple samples versus bootstrapping in coupled model. Shown are (top) time scale at 500 hPa (in days) and (*bottom*) corresponding ± 2 standard deviation, derived from (*solid*) 80 fifty-year-

long samples and (*dashed*) bootstrapping (100 times) with replacement of N = 50 year-long samples

differences between coupled and uncoupled model time scale uncertainty are even smaller. In other words, averaged over the whole atmosphere and entire year, internal atmospheric dynamics dominate the uncertainty of the time scale, while effects from low-frequency lower boundary forcing are of second order.

We next test how our uncertainty estimate derived from dividing our long control runs into shorter segments compares with an alternate method suggested by Simpson et al. (2011). Simpson's method is based on resampling one single relatively short time series using bootstrapping with replacement to produce a large number of synthetic time scale estimates. We also apply Simpson's approach, but instead of using just one arbitrary fifty-year-long model time series for the resampling, we utilize all 80 fifty-year long time series of our control run and average the results. In Fig. 9 we compare the outcomes from the two approaches for N = 50 years and a level of 500 hPa. It is reassuring to find that the two methods lead to very similar outcomes in terms of the mean time scale (top). However, resampling (dashed) consistently leads to somewhat larger uncertainty estimates than our approach (solid). Getting

back to Fig. 8, the grey symbols indicate uncertainty estimates for the NCEP/NCAR reanalysis using Simpson's method, which can be compared with the black symbols. The small differences suggests that future studies can safely use Simpson's method and that this perhaps leads to somewhat larger uncertainty estimates than our method. The huge benefit of Simpson's approach is that long time series are not needed.

8 Summary and discussion

In this study we examined the AM time scale of the GFDL general circulation model CM2.1 and compared it against the NCEP/NCAR reanalysis. We investigated the overall time-height structure of the time scale and the seasonal timing and strength of the time scale maxima in the troposphere and stratosphere. A particular focus was understanding the uncertainties associated with determining the time scale and also possible connections between stratospheric and tropospheric time scale. Overall, the model simulated time scale structure agreed reasonably well with that of the reanalysis, in particular in terms of the seasonal timing of the tropospheric peaks. However, a major model deficiency was the seasonal cycle of the time scale in the troposphere, which like in most models (Gerber et al. 2008b, 2010) was much too broad compared to the reanalysis.

An important theme of this study was the slow convergence of the time scale and the role of internal variability for the convergence. Gerber et al. (2008a) used theoretical arguments to estimate that if τ is 10 days about 10 years of a perpetual January integrations are necessary to determine the time scale at 20 % accuracy. Multiplying this number by a factor of four to account for the seasonal cycle in our integrations one arrives at about 40 years of data, which agrees very well with our empirically derived results (Fig. 8). According to our long control run at least 100 years of integration are needed in order to limit the uncertainty below 10 %. At the same time, we find that the uncertainty of the SAM time scale is consistently larger than that of the NAM, which might be indicative for larger low-frequency variability in the SAM as compared to the NAM. Our uncertainty estimates agree well with an alternate method proposed by Simpson et al. (2011), which is based on resampling. Simpson's method is preferable because it requires only short records, and, as we have shown, produces reasonable results.

The large time scale uncertainty also raises the question of whether the 50-year-long historical record is long enough to evaluate climate models. Since the uncertainty amounts to 20 % when only 50 years of data are available it is clear that caution is necessary when judging models in terms of their time scale agreement with the reanalysis. Indeed, our analysis shows that for the NAM the reanalysis lies within the scatter of outcomes from individual 50-yearlong model segments. This is not only true in terms of the overall time scale structure (Fig. 6) but also in terms of seasonal timing and magnitude of the winter maximum in NAM time scale (Fig. 7). Our results suggest that the differences between model and reanalysis can be largely explained from internal variability. Systematic model biases that would be indicative for unrealistic dynamical sensitivity could not be detected.

For the SAM, however, the above-described situation is different. For example, if the model is validated against itself under a perfect model approach, individual segments agree much better with each other than with the reanalysis. Further analysis shows that the disagreement between model and reanalysis for the SAM is mostly related to the magnitude of the spring/summer time maximum in tropospheric SAM time scale. On average and compared against the reanalysis, the model overestimates the strength of the peak in time scale by a factor of two. To the extent that this bias is associated with internal dynamics and not with overly persistent forcings, this indicates that the model's tropospheric SAM is too sensitive with respect to external influences.

There has been some recent discussion about possible influences of stratospheric variability in lengthening the tropospheric AM time scale (Baldwin et al. 2003). Various previous modeling studies found conclusive evidence for such influences. We also investigated this question and found that in the coupled model the stratospheric connection to the tropospheric time scale is weak at best over the NH and quite strong over the SH (Figs. 6, 7). The situation in the uncoupled model is somewhat different, with a weaker connection over the SH and a stronger connection over the NH. The correlation over the NH is still quite weak and therefore somewhat in disagreement with Simpson et al. (2011), who state that "in the NH, virtually all of the seasonality in AM time scales arises from the downward influence of stratospheric annular mode variability". It is likely that differences in how the two studies approach the question and how stratospheric variability is generated are responsible for the different conclusions.

The connection between stratospheric and tropospheric time scale in the model is quite pronounced over the SH in terms of the magnitude of the peak time scale (Fig. 7d). For the SAM, there is only a slight delay between stratospheric and tropospheric time scale maxima, and most model segments suggest that over the SH the stratospheric peak leads the troposphere by a week or so. Further, there are no significant differences between model and reanalysis in terms of this lag over the SH. For the NH, the situation is very different. Averaged over all model segments, the stratospheric peak in the NAM leads the tropospheric peak by about 2 months, which is unrealistically long when compared to the reanalysis (Figs. 2, 7). However, as noted above, many individual model segments are in perfect agreement with the reanalysis in terms of their timing of stratospheric and tropospheric peaks. Based on these results one must conclude that differences in the stratosphere-troposphere lag between the mean model outcome and the reanalysis are not indicative for real model deficiencies, or that the lag of just a few days seen in the reanalysis is a characteristic feature of our atmosphere.

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References

- Ambaum MHP, Hoskins BJ (2002) The NAO troposphere-stratosphere connection. J Clim 15:1969–1978
- Baldwin MP, Thompson DWJ (2009) A critical comparison of stratosphere-troposphere coupling indices. Quart J R Met Soc 135:1661–1672
- Baldwin MP, Stephenson DB, Thompson DWJ, Dunkerton TJ, Charlton AJ, Alan ON (2003) Stratospheric memory and skill of extended-range weather forecasts. Science 301:636–640
- Barsugli JJ, Battisti DS (1998) The basic effects of atmosphereocean thermal coupling on midlatitude variability. J Atmos Sci 55:477–493
- Chan CJ, Plumb RA (2009) The response to stratospheric forcing and its dependence on the state of the troposphere. J Atmos Sci 66:2107–2115
- Chen G, Plumb RA (2009) Quantifying the eddy feedback and the persistence of the zonal index in an idealized atmospheric model. J Atmos Sci 66:3707–3720
- Delworth TL et al (2006) GFDL's CM2 global coupled climate models. Part I: formulation and simulation characteristics. J Clim 19:643–674
- Gerber EP, Polvani LM (2009) Stratosphere-troposphere coupling in a relatively simple AGCM: the importance of stratospheric variability. J Clim 22:1920–1933
- Gerber EP, Voronin S, Polvani LM (2008a) Testing the annular mode autocorrelation time scale in simple atmospheric general circulation models. Mon Wea Rev 136:1523–1536
- Gerber EP, Polvani LM, Ancukiewicz D (2008b) Annular mode time scales in the Intergovernmental Panel on Climate Change Fourth Assessment Report models. Geophys Res Lett 35:L22707
- Gerber EP et al (2010) Stratosphere-troposphere coupling and annular mode variability in chemistry-climate models. J Geophys Res 115:D00M06
- GFDL Global Atmospheric Model Development Team (2004) The new GFDL global atmosphere and land model AM2-LM2: evaluation with prescribed SST simulations. J Clim 17:4641–4673
- Grise KM, Polvani LM (2014) Is climate sensitivity related to dynamical sensitivity? A Southern Hemisphere perspective. Geophys Res Lett 41:534–540

- Houghton JT et al (2001) Climate change 2001: The scientific basis. Cambridge University Press, Cambridge
- Kalnay E et al (1996) The NCEP/NCAR 40-year reanalysis project. Bull Am Met Soc 77:437–471
- Keeley SPE, Sutton RT, Shaffrey LC (2009) Does the North Atlantic oscillation show unusual persistence on intraseasonal timescales? Geophys Res Lett 36:L22706
- Kidston J, Gerber EP (2010) Intermodel variability of the poleward shift of the austral jet stream in the CMIP3 integrations linked to biases in 20th century climatology. Geophys Res Lett 37:L09708
- Knutti R, Meehl GA, Allen MR, Stainforth DA (2006) Constraining climate sensitivity from the seasonal cycle in surface temperature. J Clim 19:4224–4233
- Leith CE (1975) Climate response and fluctuation dissipation. J Atmos Sci 32:2022–2026
- Meehl GA et al (2007) Global Climate Projections. In: Solomon S, Qin D, Manning M, Chen Z, Marquis M, Averyt KB, Tignor M, Miller HL (eds) Climate Change 2007: The physical science basis. Contribution of Working Group I to the fourth assessment report of the intergovernmental panel on climate change. Cambridge University Press, Cambridge
- Randall DA et al (2007) Cilmate models and their evaluation. In: Solomon S, Qin D, Manning M, Chen Z, Marquis M, Averyt KB, Tignor M, Miller HL (eds) Climate change 2007: The physical science basis. Contribution of working Group I to the fourth assessment report of the intergovernmental panel on climate change. Cambridge University Press, Cambridge
- Reichler T, Kim J (2008) How well do coupled models simulate today's climate? Bull Am Met Soc 89:303–311
- Reichler T, Kim J, Manzini E, Kröger J (2012) A stratospheric connection to Atlantic climate variability. Nature Geos 5:783–787
- Ring MJ, Plumb RA (2008) The response of a simplified GCM to axisymmetric forcings: applicability of the fluctuation–dissipation theorem. J Atmos Sci 65:3880–3898
- Schimanke S, Körper J, Spangehl T, Cubasch U (2011) Multi-decadal variability of sudden stratospheric warmings in an AOGCM. Geophys Res Lett 38:L01801
- Simpson IR, Hitchcock P, Shepherd TG, Scinocca JF (2011) Stratospheric variability and tropospheric annular-mode timescales. Geophys Res Lett 38:L20806
- Simpson IR, Hitchcock P, Shepherd TG, Scinocca JF (2013a) Southern annular mode dynamics in observations and models. Part I: the influence of climatological zonal wind biases in a comprehensive GCM. J Clim 26:3953–3967
- Simpson IR, Shepherd TG, Hitchcock P, Scinocca JF (2013b) Southern annular mode dynamics in observations and models. Part II: Eddy Feedbacks. J Clim 26:5220–5241
- Son S-W, Lee S, Feldstein SB, Ten Hoeve JE (2008) Time scale and feedback of zonal-mean-flow variability. J Atmos Sci 65:935–952
- Staten PW, Reichler T, Lu J (2014) The transient circulation response to radiative forcings and sea surface warming. J Clim 27:9323–9336
- Thompson DWJ, Wallace JM (2000) Annular modes in the extratropical circulation. Part I: Month-to-month variability. J Clim 13:1000–1016
- Thompson DWJ, Wallace JM (2001) Regional climate impacts of the Northern Hemisphere annular mode. Science 293:85–89