

# QCOM 2D Convection

$$\frac{\partial v}{\partial t} = -v \frac{\partial v}{\partial y} - w \frac{\partial v}{\partial z} - c_p \theta_0 \frac{\partial \pi_1}{\partial y} + D_v,$$

$$\frac{\partial w}{\partial t} = -v \frac{\partial w}{\partial y} - w \frac{\partial w}{\partial z} - c_p \theta_0 \frac{\partial \pi_1}{\partial z} + g \left( \frac{\theta}{\theta_0} - 1 \right) + D_w,$$

$$\frac{\partial \theta}{\partial t} = -v \frac{\partial \theta}{\partial y} - w \frac{\partial \theta}{\partial z} + D_\theta,$$

$$\frac{\partial \pi_1}{\partial t} = -\frac{c_s^2}{c_p \theta_0^2} \left[ \frac{\partial}{\partial y} (\theta_0 v) + \frac{\partial}{\partial z} (\theta_0 w) \right].$$

The right-hand sides of (1)-(4) will be denoted  $f_v$ ,  $f_w$ ,  $f_\theta$ , and  $f_\pi$ . Their centered, second-order accurate finite-difference forms are

$$f_v = -v\delta_{2y}v - \overline{\bar{w}^y\delta_z v^z} - c_p\theta_0\delta_y\pi_1 + D_v, \quad (6)$$

$$f_w = -\overline{\bar{v}^z\delta_y w^y} - w\delta_{2z}w - c_p\overline{\bar{\theta}_0^z}\delta_z\pi_1 + g(\frac{\bar{\theta}^z}{\overline{\bar{\theta}_0^z}} - 1) + D_w, \quad (7)$$

$$f_\theta = -\overline{\bar{v}\delta_y\theta^y} - \overline{\bar{w}\delta_z\theta^z} + D_\theta, \quad (8)$$

$$f_\pi = -\frac{c_s^2}{c_p\theta_0^2} \left[ \delta_y(\theta_0 v) + \delta_z(\overline{\bar{\theta}_0^z} w) \right]. \quad (9)$$

Each of (1)-(4) can be written as

$$\frac{\partial\phi}{\partial t} = f_\phi,$$

## 4 Turbulence closure

For simplicity, we will use the eddy viscosity approach. Then terms  $D_v$ ,  $D_w$ , and  $D_\theta$  each have the form

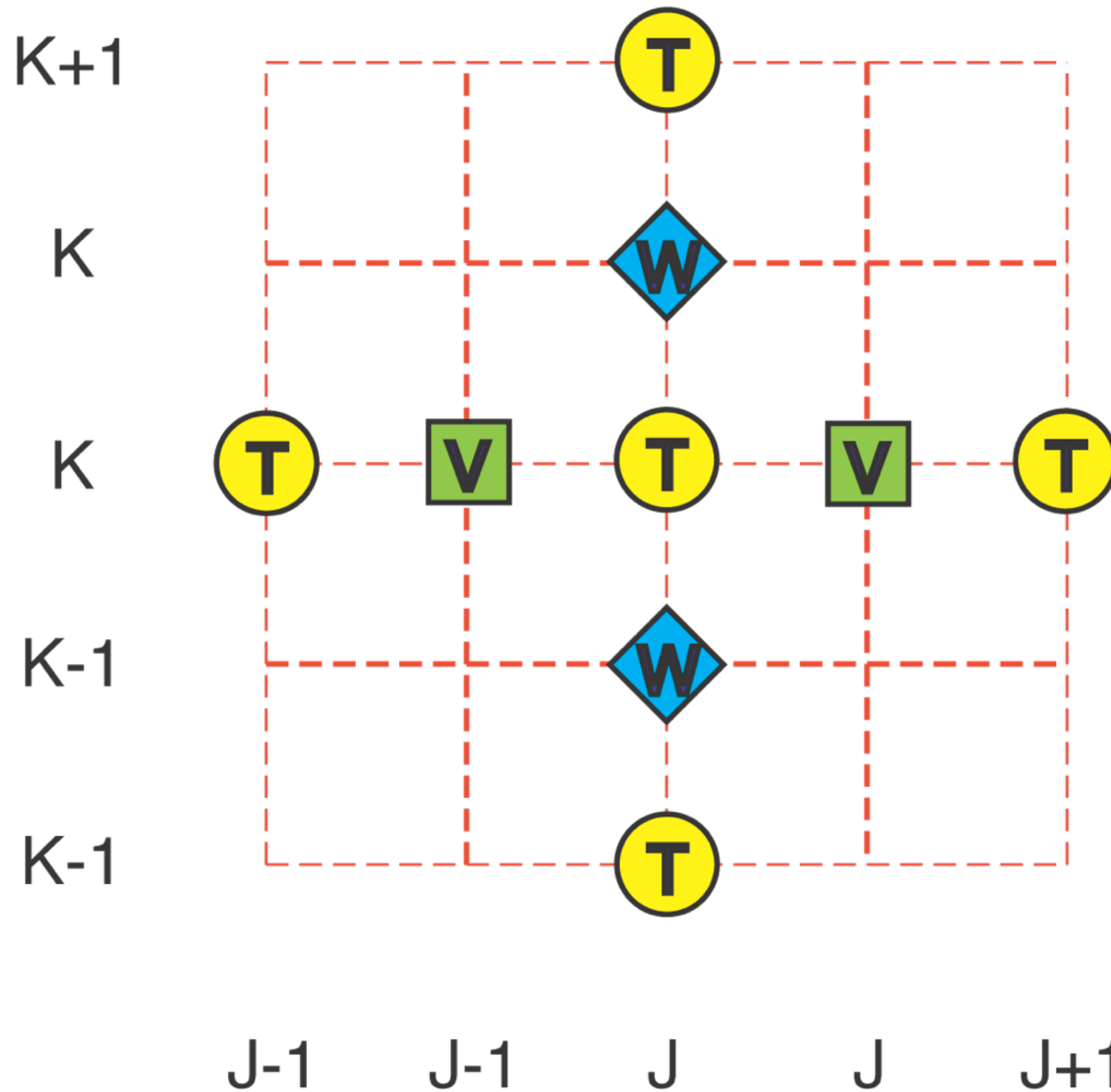
$$K_\phi \nabla^2 \phi,$$

or, in finite-difference form,

$$K_\phi [\delta_y(\delta_y \phi) + \delta_z(\delta_z \phi)],$$

where  $K_\phi$  is the eddy diffusivity.

# Array subscripting convention



## Analytic Solution to the Linear Convection Equations

